

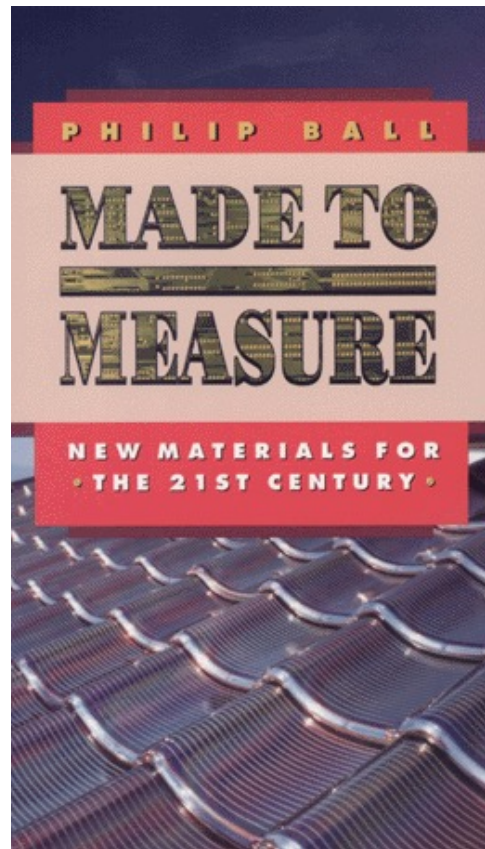


Non-equilibrium interfaces

Margarida Telo da Gama

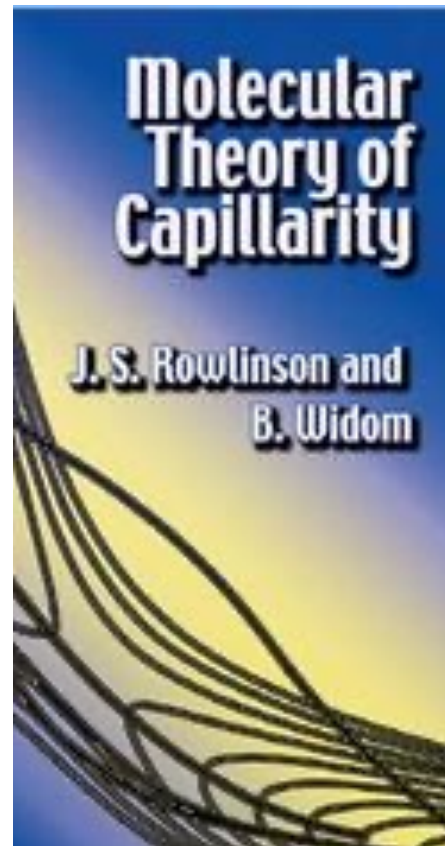
There seems to exist a certain degree of prejudice against surfaces, with things superficial being deemed of lesser pith and moment.

P Ball, 1997



This is most unjustified as
'surfaces and interfaces are the
seat of a variety of fascinating
phenomena'.

J S Rowlinson and B Widom,
1982

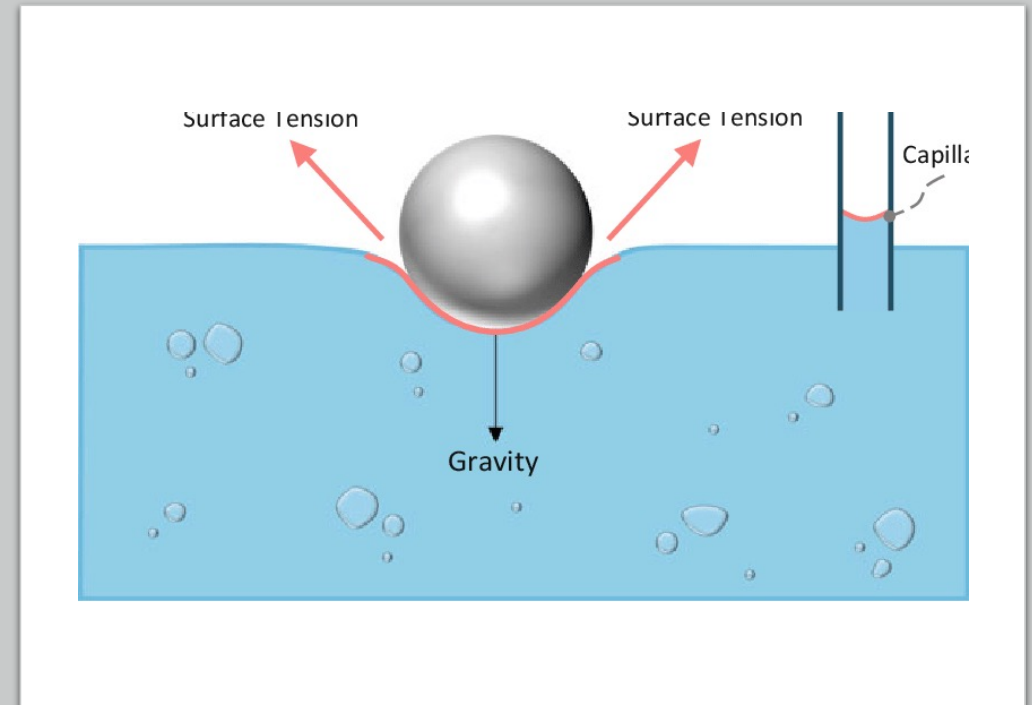
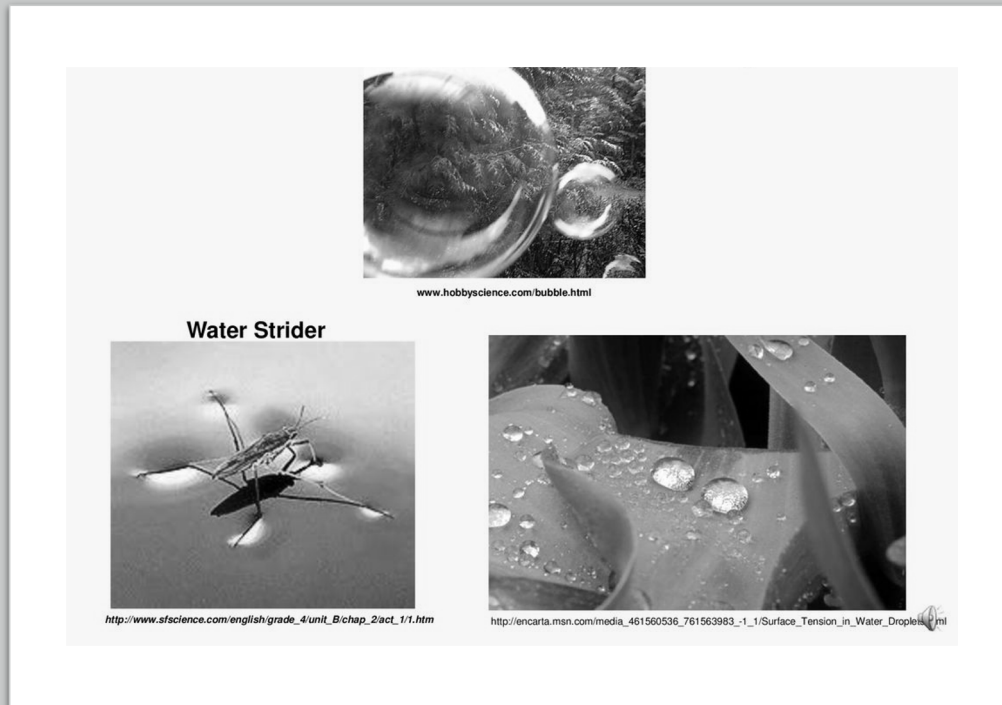




Equilibrium interfaces

Surface tension

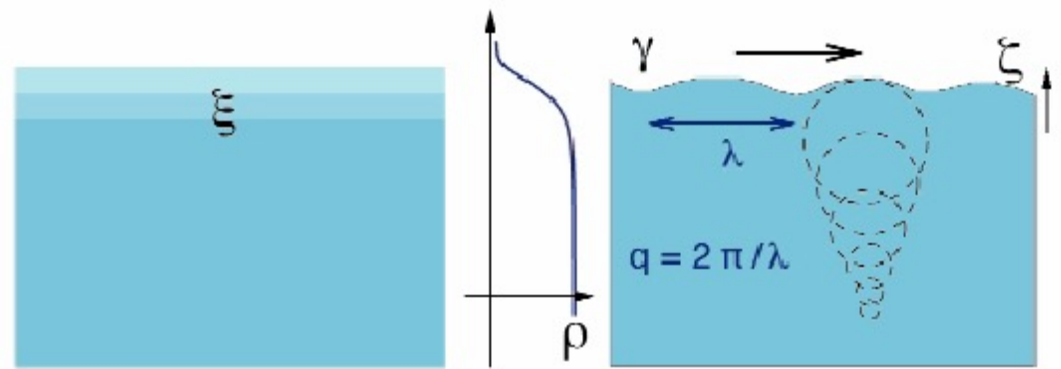
- Excess free energy per unit area
- Tangential force per unit length



The balance between the cohesion of the liquid and its adhesion to the material of the container determines the degree of wetting, the contact angle and the shape of meniscus.

van der Waals and CW model

- The van der Waals description of a liquid-vapour interface (1893) predicts a smooth transition from the liquid density to the vapour density with a characteristic length related to the bulk correlation length.
- In the capillary wave model (1965), the interface structure is due to the propagation of thermally excited capillary waves with wave-vector q . Their amplitude is inversely proportional to the surface tension γ .



$$\zeta_q^2 = \frac{2kT}{L\gamma q^2}$$

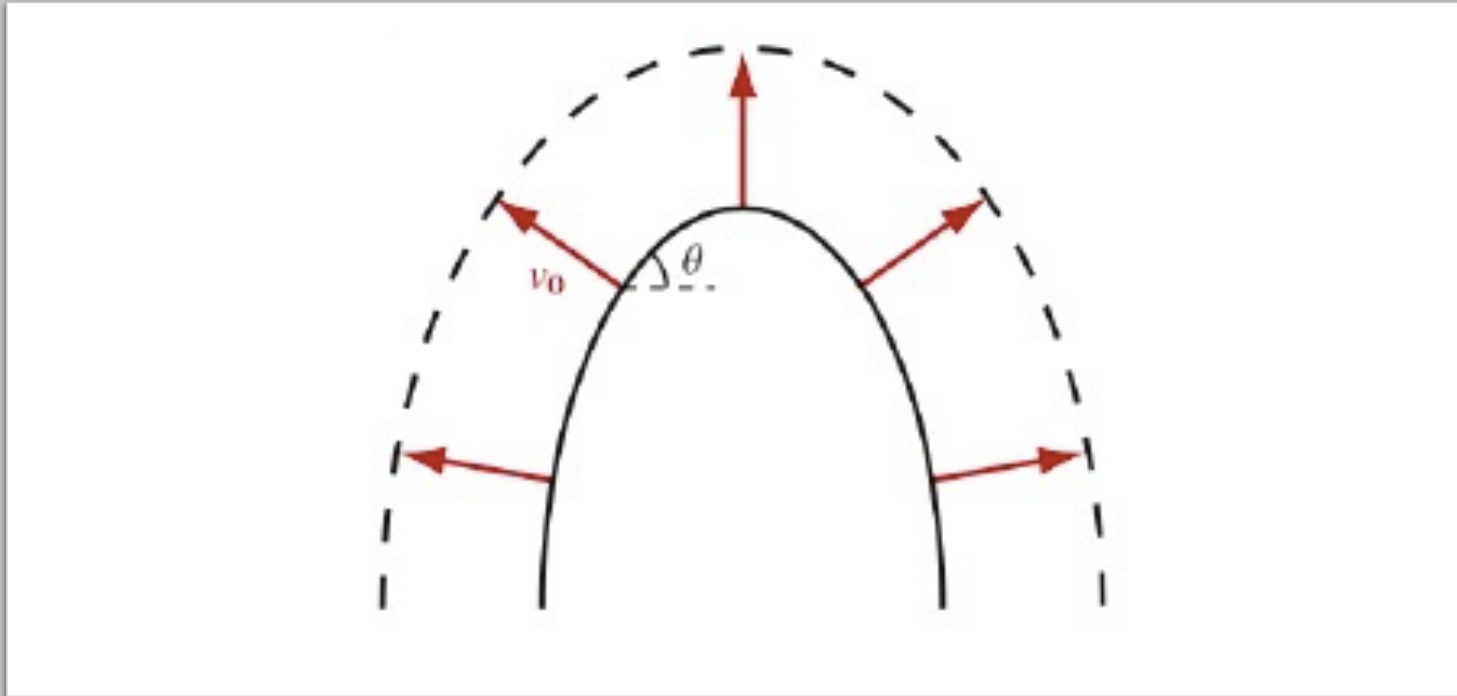
F. P. Buff, R. A. Lovett, and F. H. Stillinger, Jr., Phys. Rev. Lett. **15**, 621 (1965)



Wet or non-wet ?
Rough or smooth ?
Wetting or roughening transitions ?



Non-equilibrium wetting

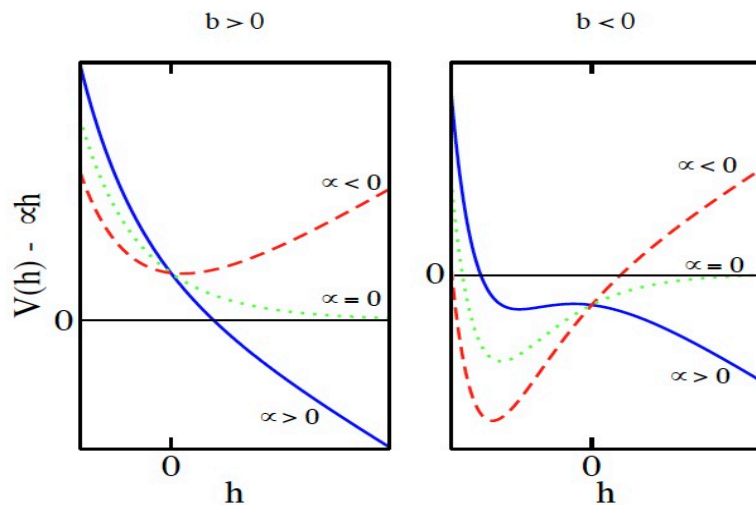


Generic interfacial growth (KPZ for $\lambda \neq 0$)

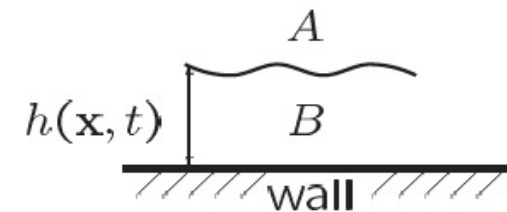
$$\frac{\partial h(x, t)}{\partial t} = \nu \frac{\partial^2 h(x, t)}{\partial x^2} + \lambda \left(\frac{\partial h(x, t)}{\partial x} \right)^2 + \eta(x, t),$$

Bound AB interface

Phases A and B coexist when $\mu = 0$

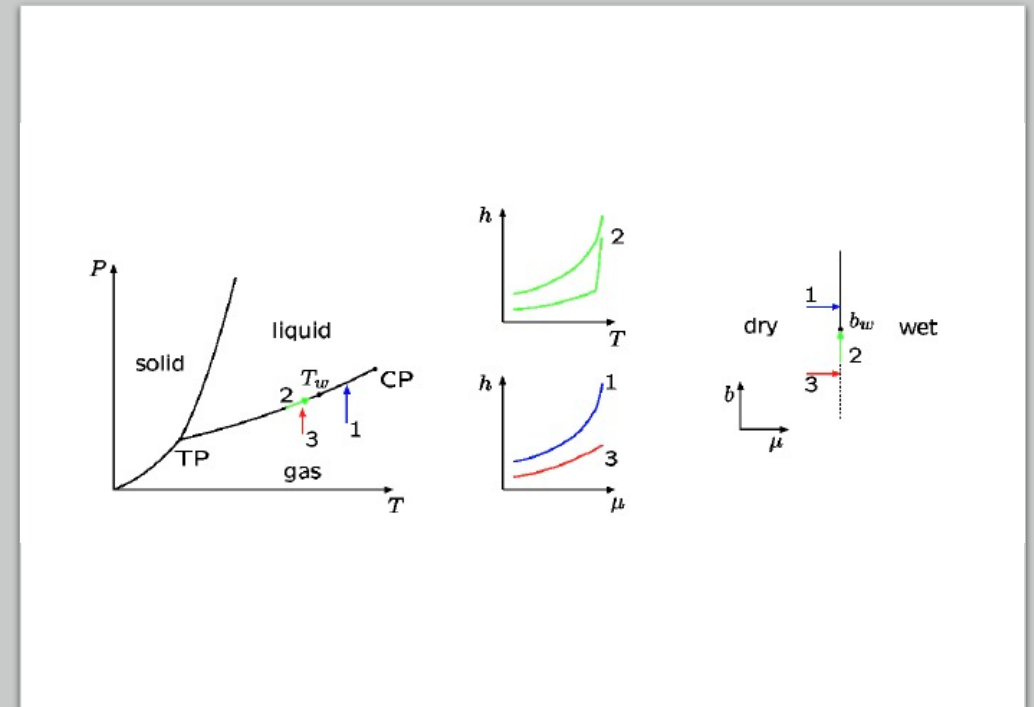
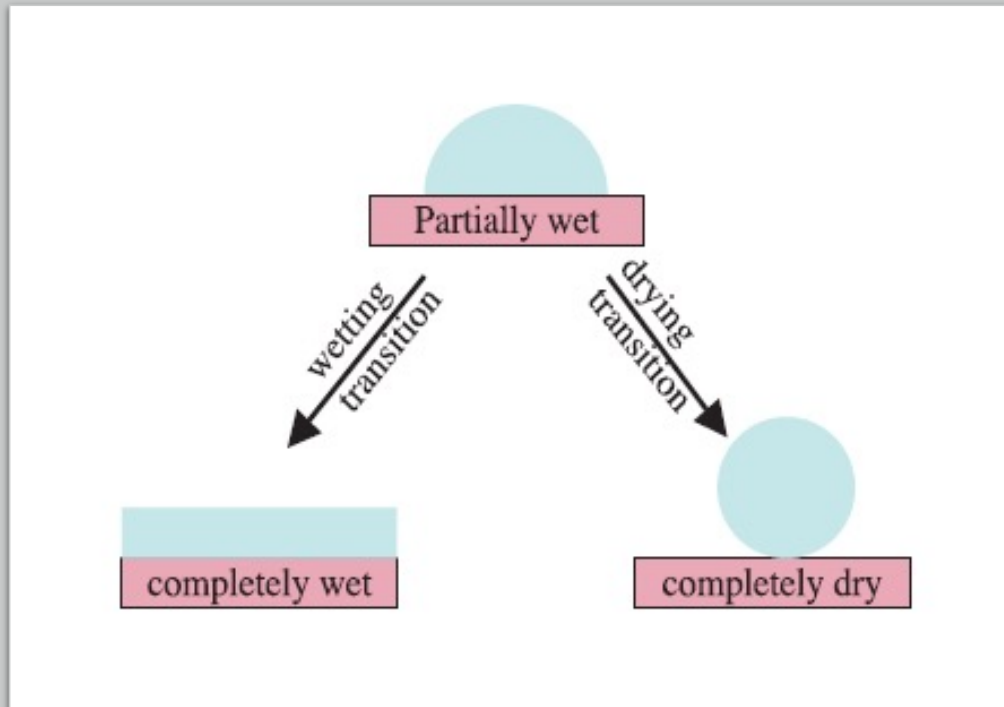


$$V(h) = b(T)e^{-h} + ce^{-2h},$$



$$\frac{\partial h(\mathbf{x}, t)}{\partial t} = \nu \nabla^2 h + \mu - \frac{\partial V}{\partial h} + \eta(\mathbf{x}, t),$$

Equilibrium wetting transition

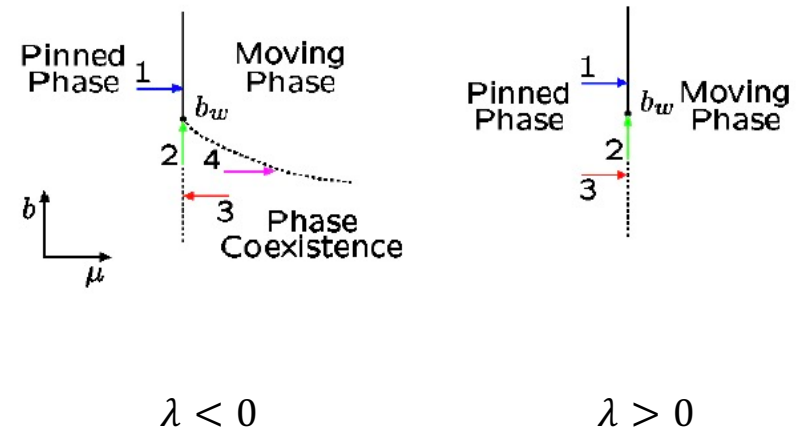


Non-equilibrium wetting & depinning

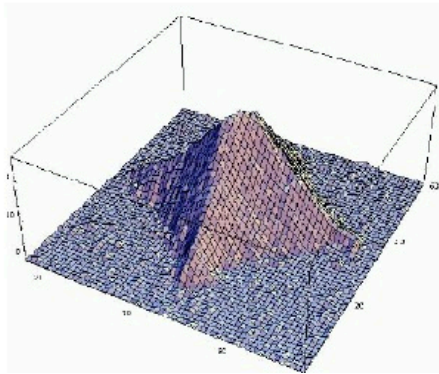
$$\frac{\partial h(\mathbf{x}, t)}{\partial t} = \nu \nabla^2 h + \lambda (\nabla h)^2 + \mu - \frac{\partial V(h)}{\partial h} + \eta(\mathbf{x}, t).$$

Wetting and drying are NOT symmetric: there is an extended coexistence for depinning & not for pinning

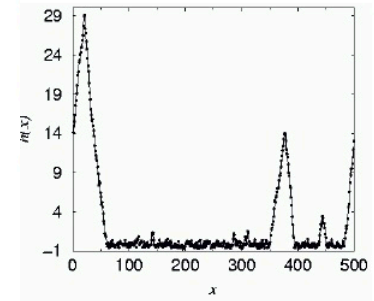
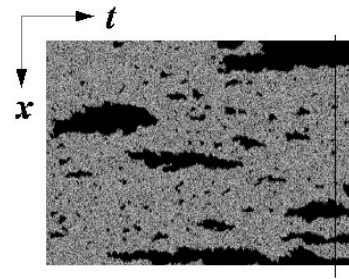
Different exponents



Extended coexistence of pinned and depinned phases in 2d & 1d



$$\lambda < 0$$



$$\lambda < 0$$

Thanks

Francisco de los
Santos

Miguel Angel Muñoz

Stochastic Theory of Non-Equilibrium Wetting, F. de los Santos, M. M. Telo da Gama and M. A. Muñoz, *Europhysics Letters*, **57**, 803-809 (2002).

Non equilibrium Wetting Transitions with Short Range Forces, F. de los Santos, M. M. Telo da Gama and M. A. Muñoz, *Physical Review E*, **67**, 021607 (2003).

Nonequilibrium Bound Interfaces, F. de los Santos e M. M. Telo da Gama in *Trends in Statistical Physics*, (World Scientific, 2004).



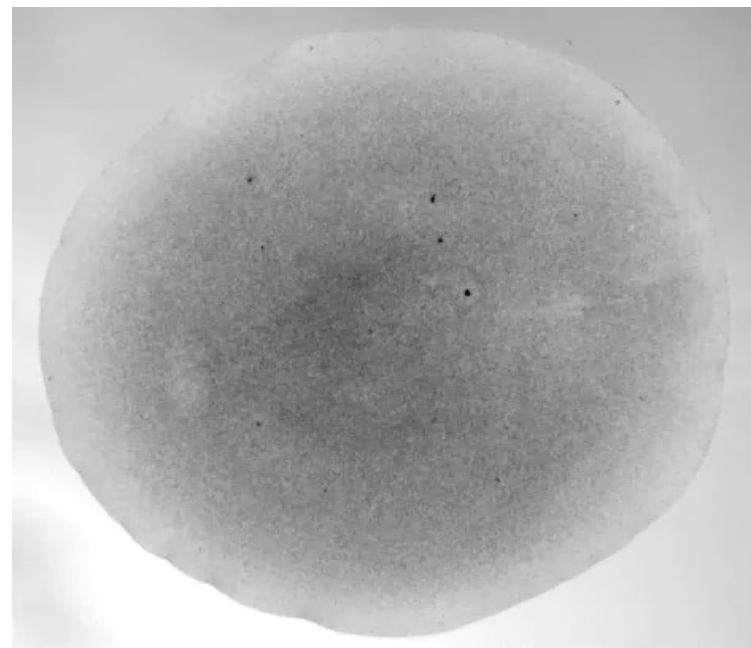
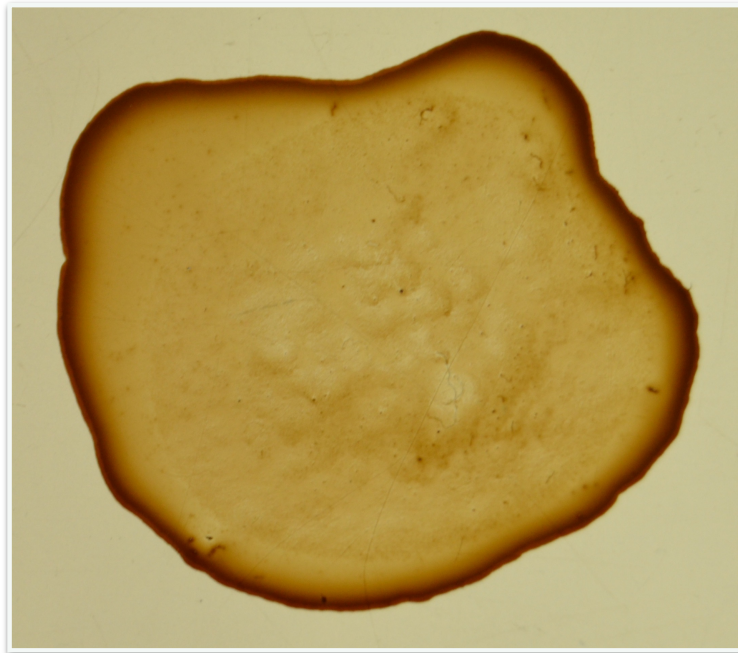


Interfacial roughening of anisotropic colloids

Evaporating drop

Coffee-ring left by spherical colloidal particles

Inhomogeneous density

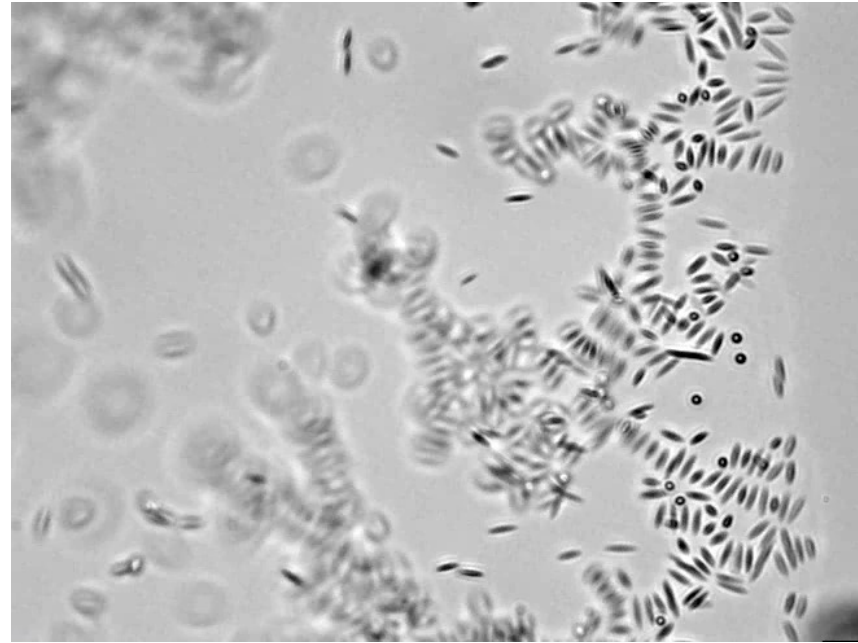
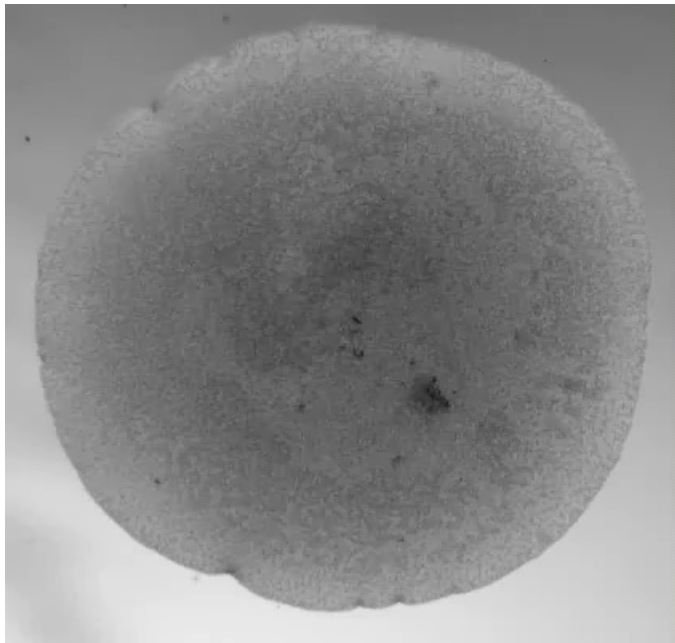


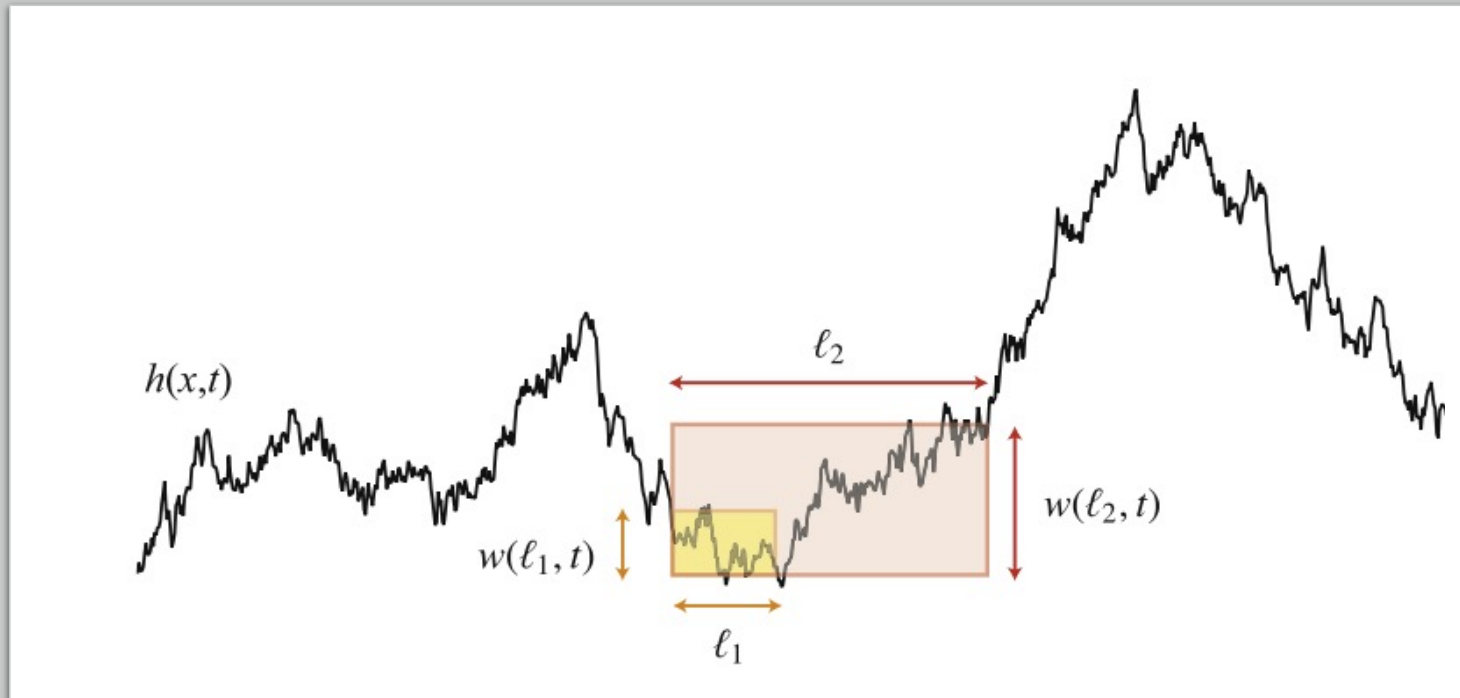
P.J. Yunker, T. Still, M.A. Lohr, A.G. Yodh, *Nature* **476**, 308 (2011)

Evaporating drop

No coffee-ring left by anisotropic colloidal particles

Homogeneous density



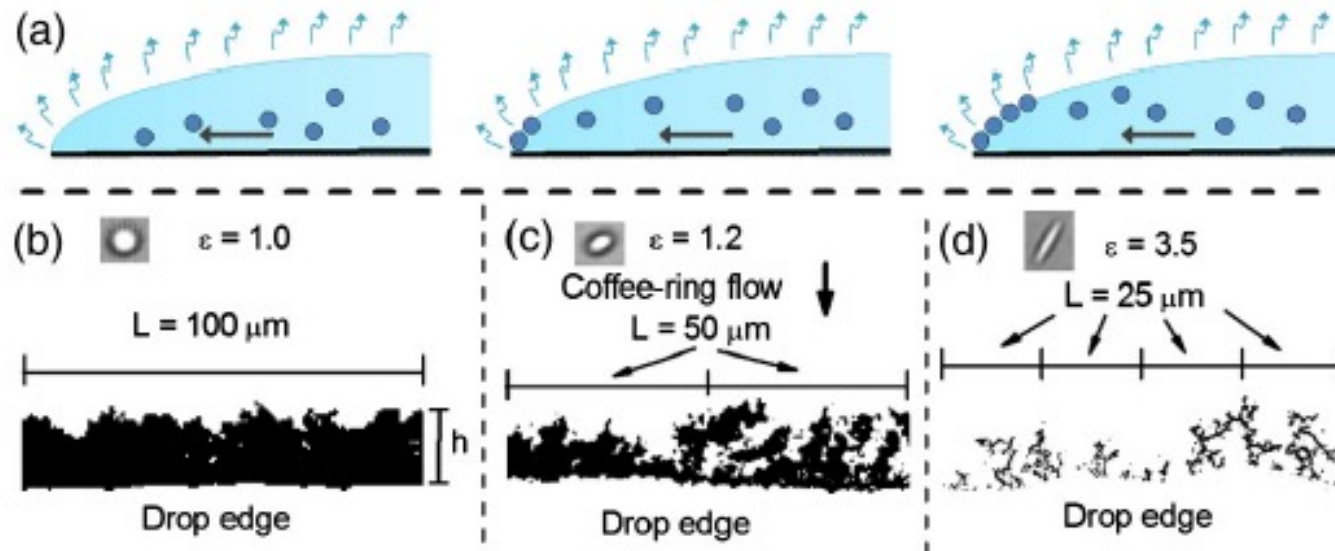


Interfacial width

$$w(t) = \sqrt{\langle [h_i(t) - \langle h(t) \rangle]^2 \rangle},$$

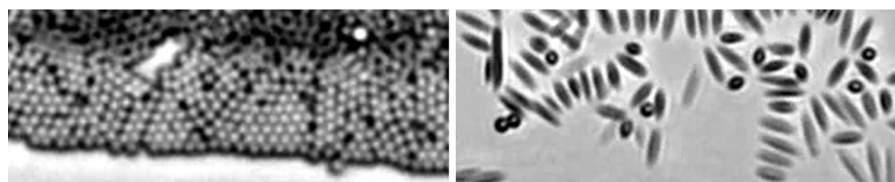
$$w(L, t) = L^\alpha f\left(\frac{t}{L^z}\right),$$

Different interfacial roughness regimes



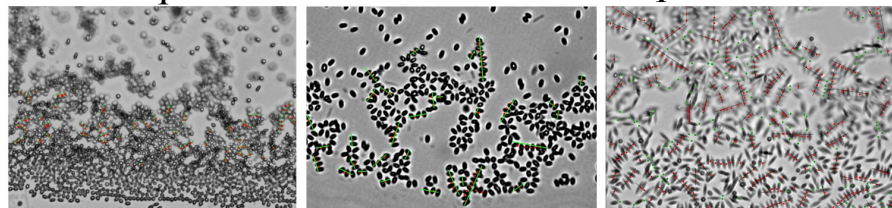
Transition from KPZ to KPZQ for anisotropic colloids

Chains & junctions driven by capillarity



Spheres

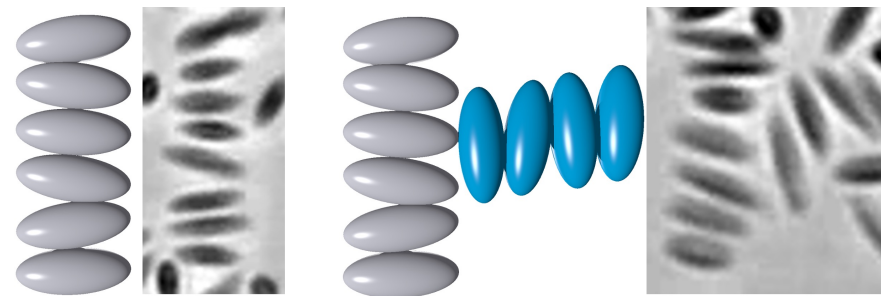
Ellipsoids



$\epsilon=1.2$

$\epsilon=1.5$

$\epsilon=3.5$

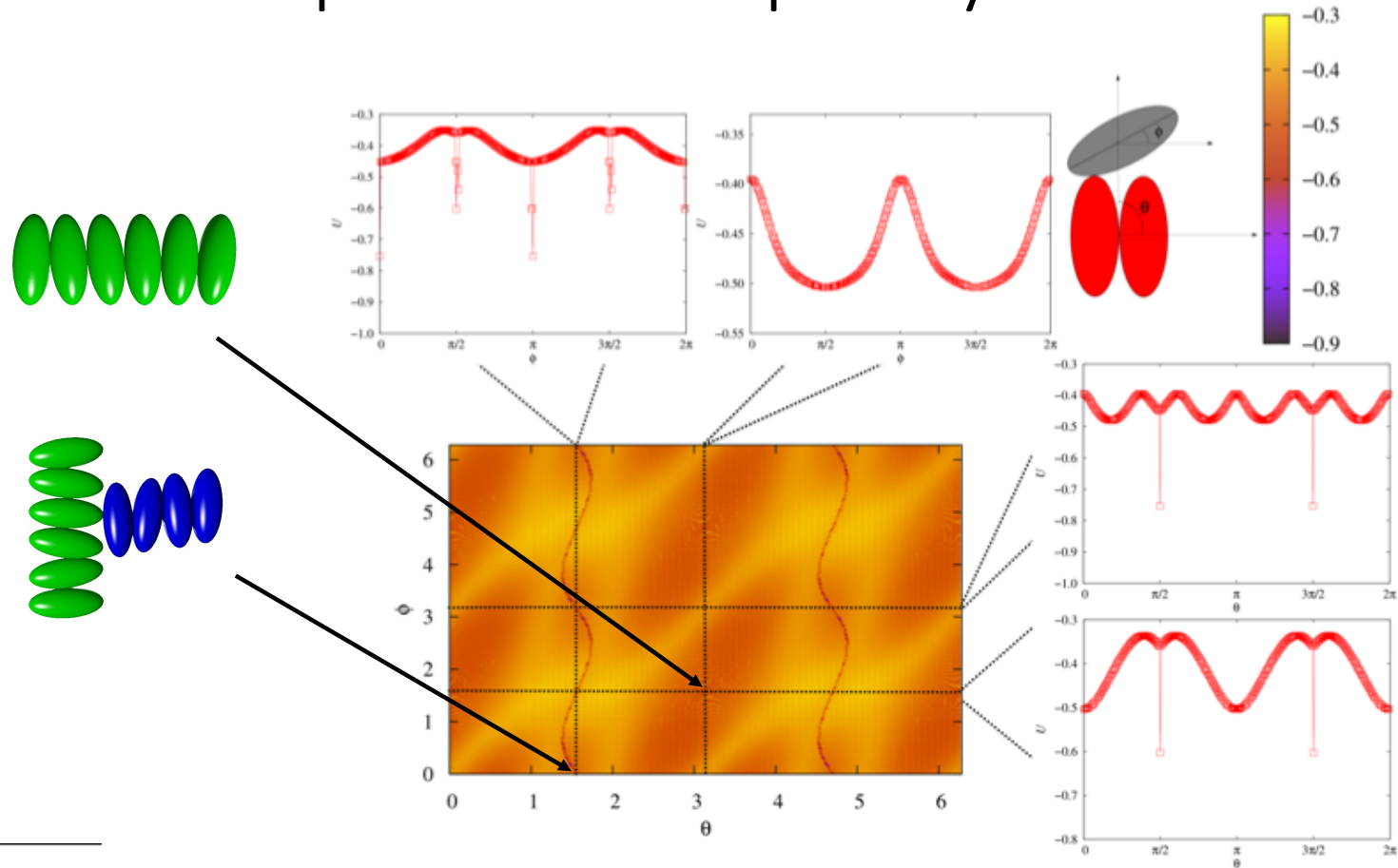


Side-to-side (chain)

Pole-to-side (junction)

C.S. Dias, P.J. Yunker, A.G. Yodh, N.A.M. Araújo, M.M. Telo da Gama, *Soft Matter* **14**, 1903 (2018)

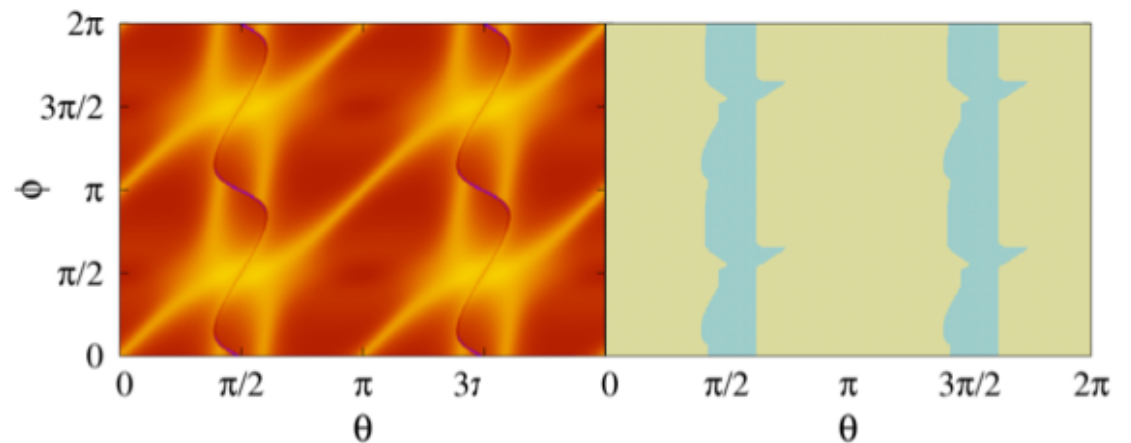
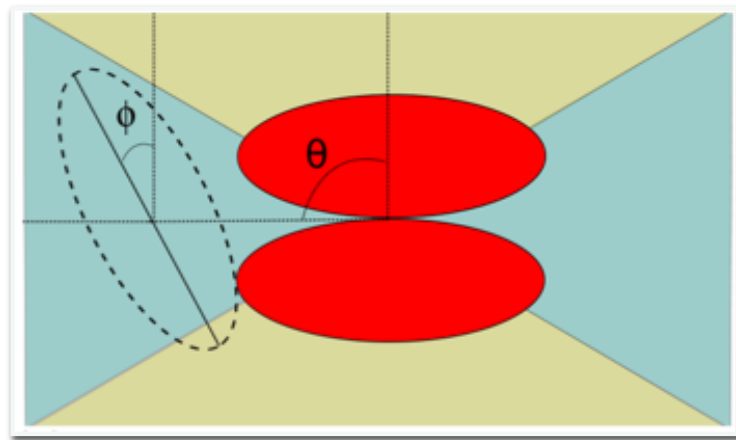
Energy landscape of the capillary interactions



Derjaguin

$$U \approx -\frac{\pi^2}{24} \sqrt{\frac{2}{h_0 (C_1 + C_2)}}$$

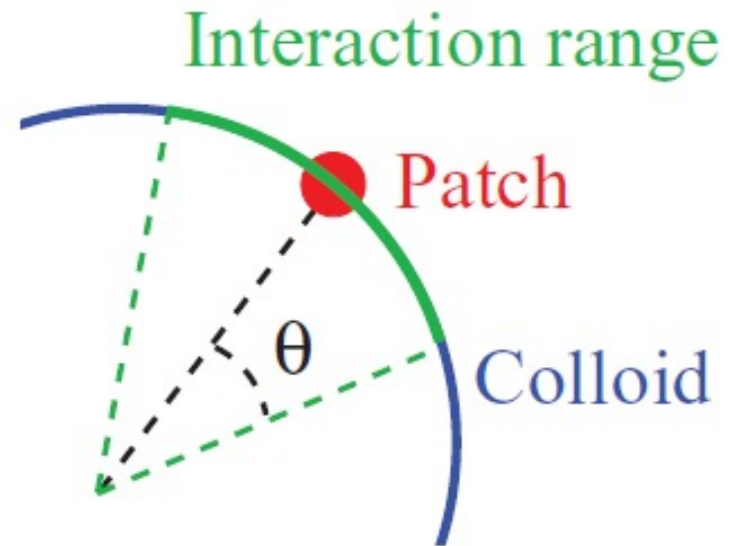
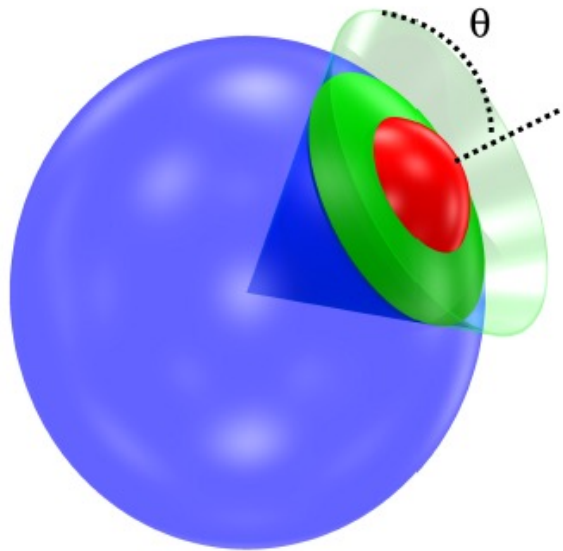
From the basins of attraction to bonding probabilities P_{AB}



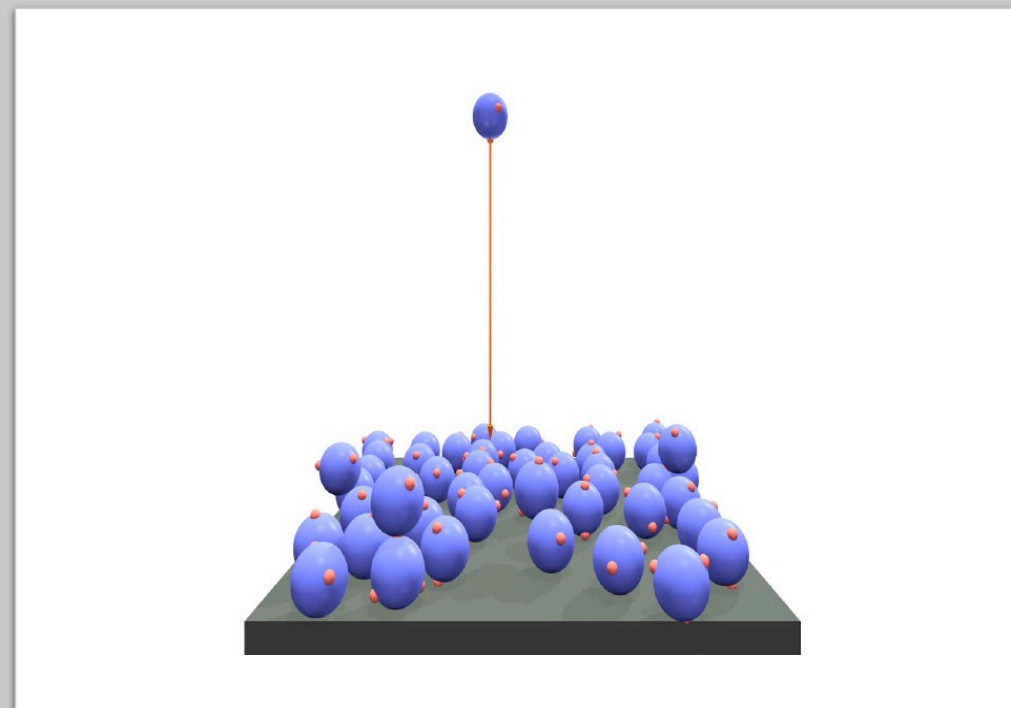
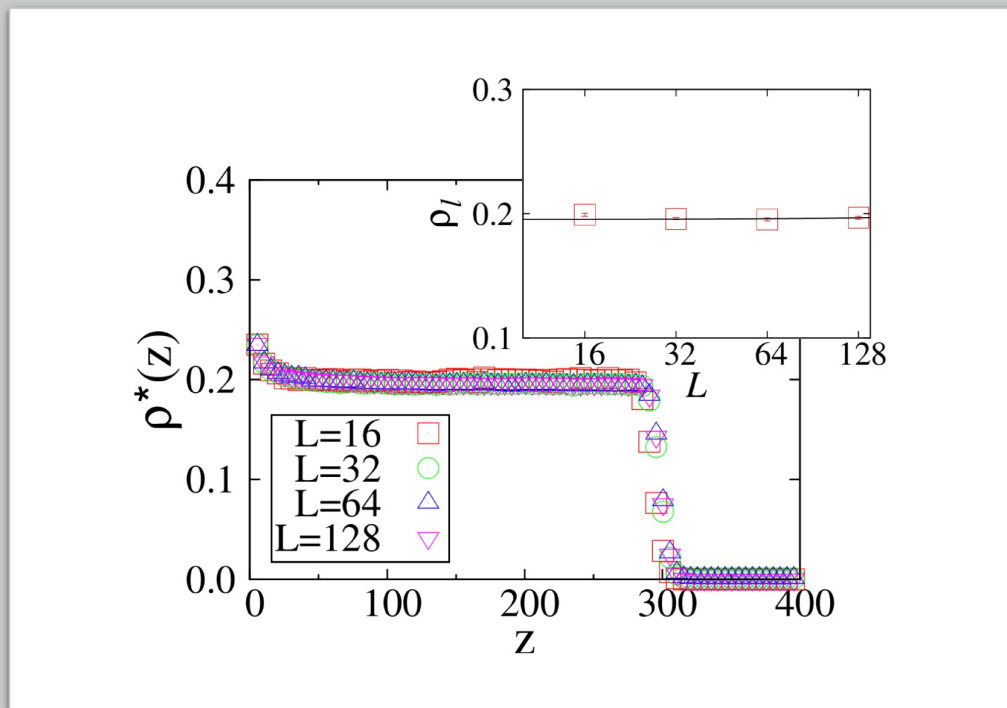
C.S. Dias, P.J. Yunker, A.G. Yodh, N.A.M. Araújo, M.M. Telo da Gama, *Soft Matter* **14**, 1903 (2018)

K.J. Schrenk, N.A.M. Araújo, J.S. Andrade Jr., H.J. Herrmann, *Scientific Reports* **2**, 348 (2012)

Stochastic kinetic model for patchy particles: Binding probability P_{AB}

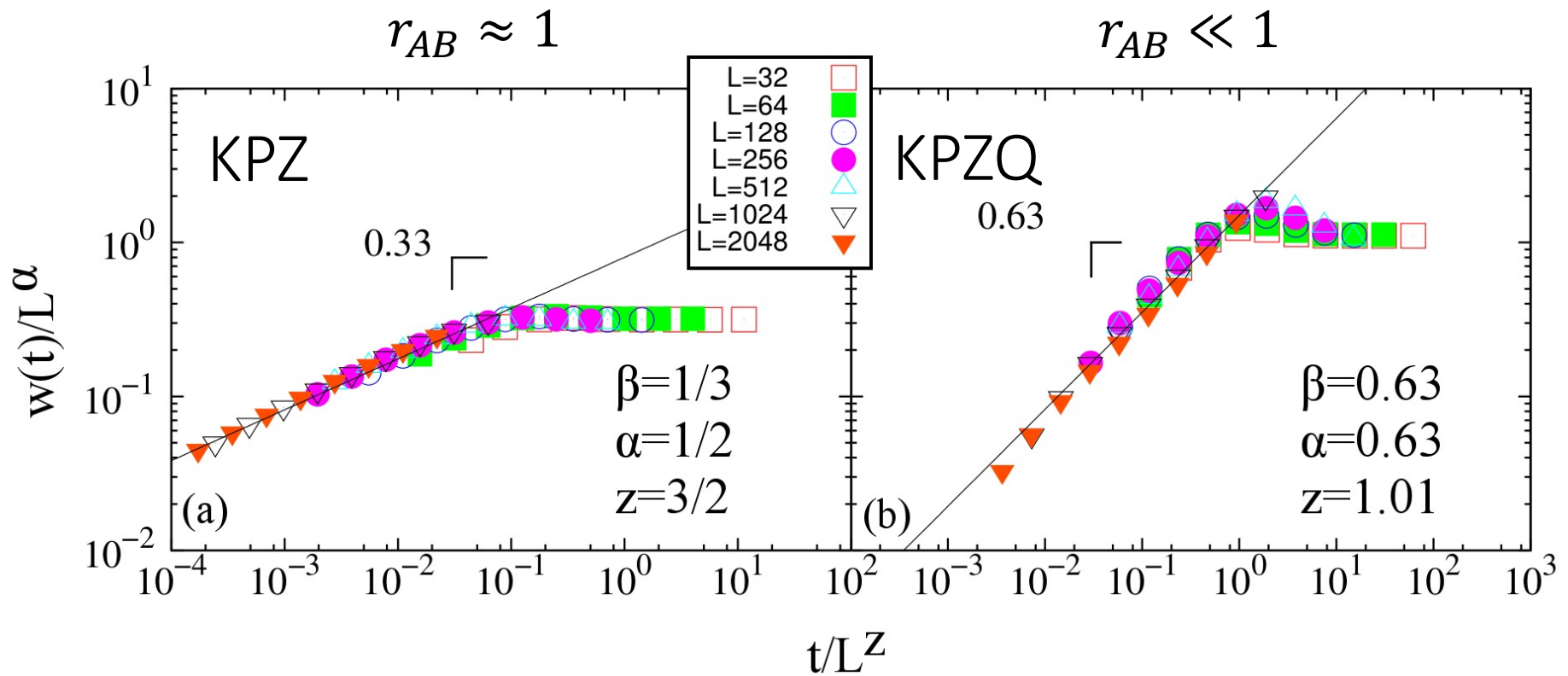


Mass transport: advection



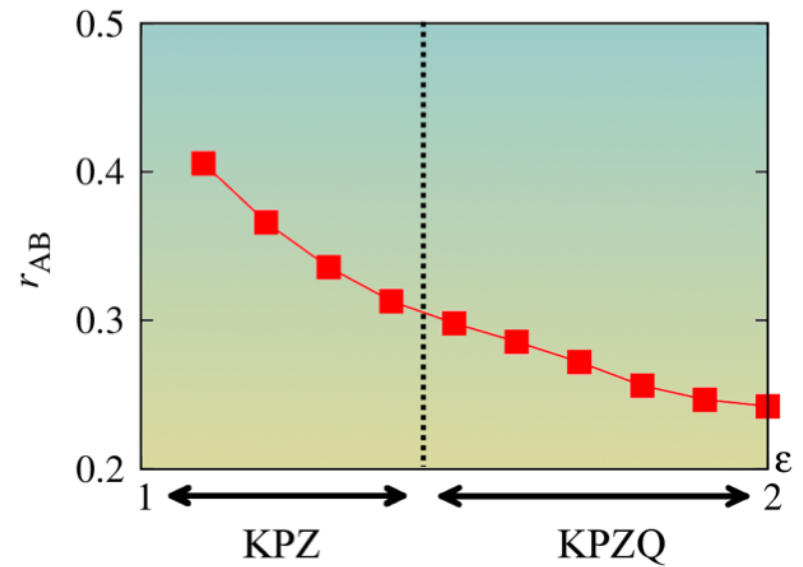
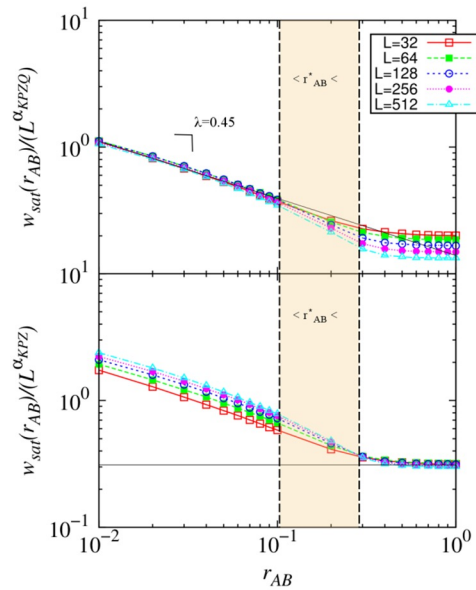
Anisotropic bonding

$$r_{AB} = \frac{P_{AB}}{P_{AA}}$$



Interfacial width

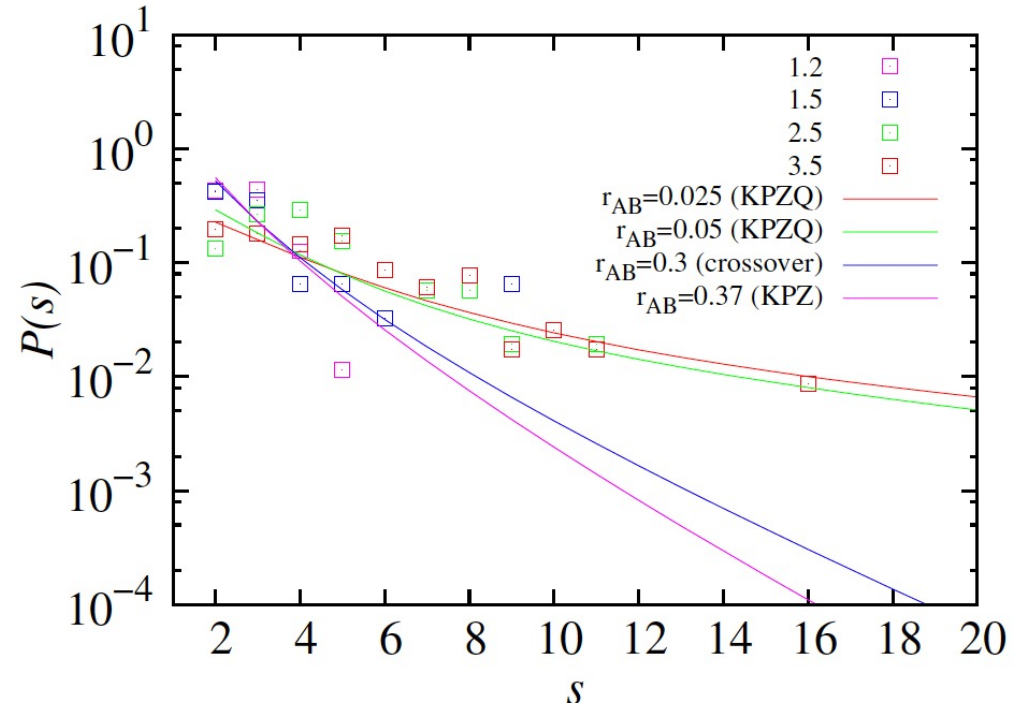
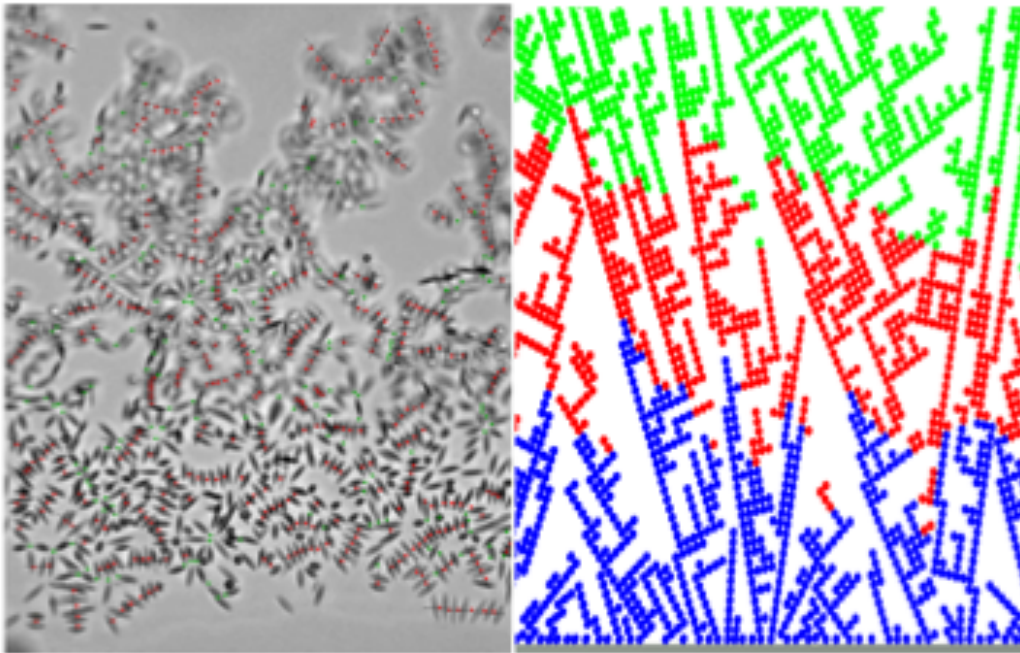
$$r_{AB} = \frac{P_{AB}}{P_{AA}}$$



C. S. Dias, N. A. M. Araújo, and M. M. Telo da Gama. EPL **107**, 56002 (2014).

C.S. Dias, P.J. Yunker, A.G. Yodh, N.A.M. Araújo, M.M. Telo da Gama, *Soft Matter* **14**, 1903 (2018)

Cluster size distributions

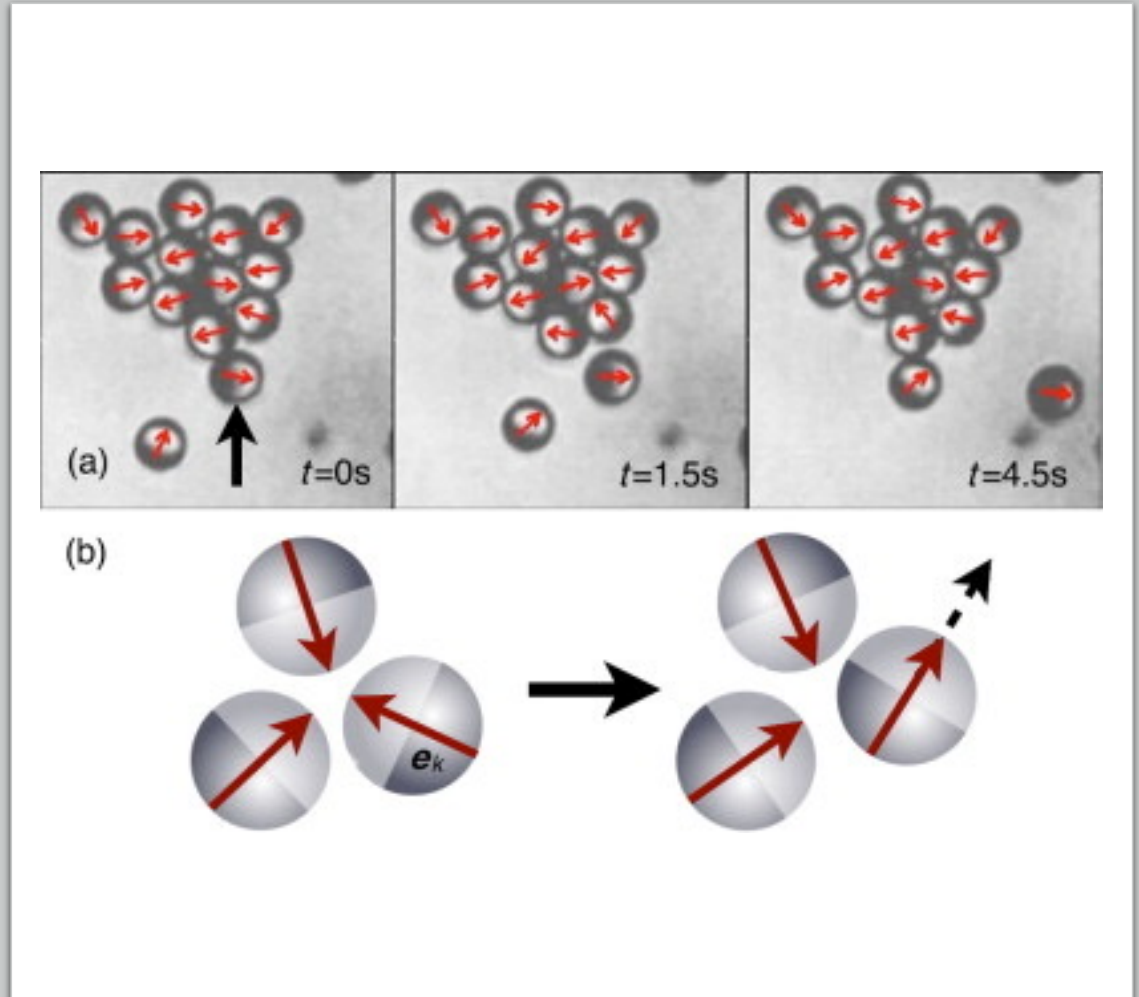




Active versus driven particles

New phases

Motility Induced Phase Separation

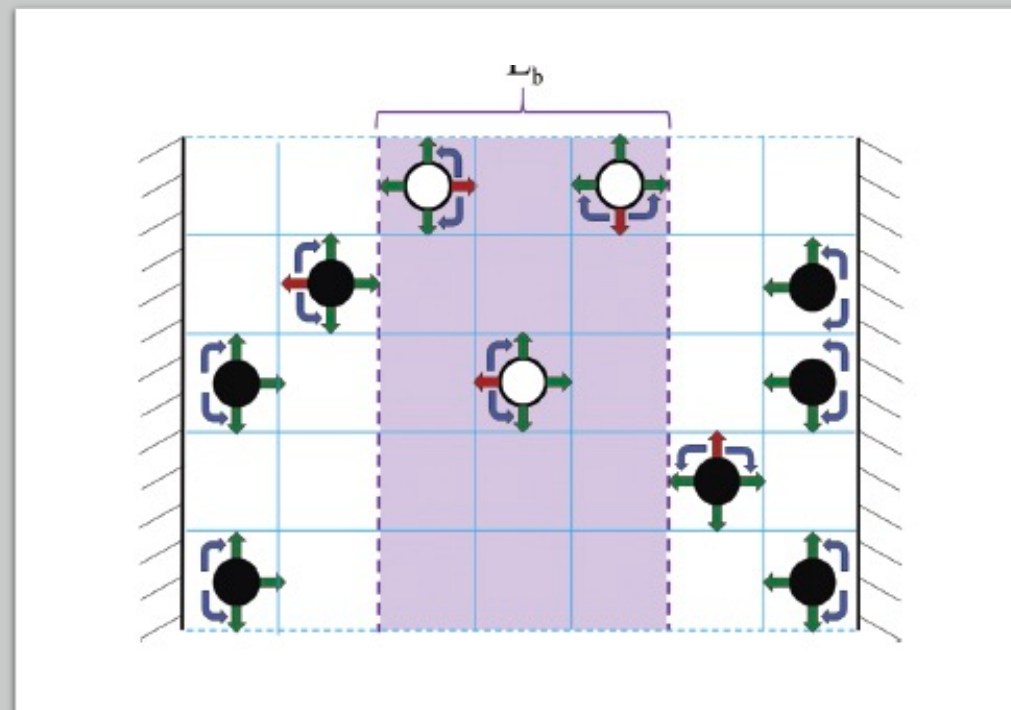
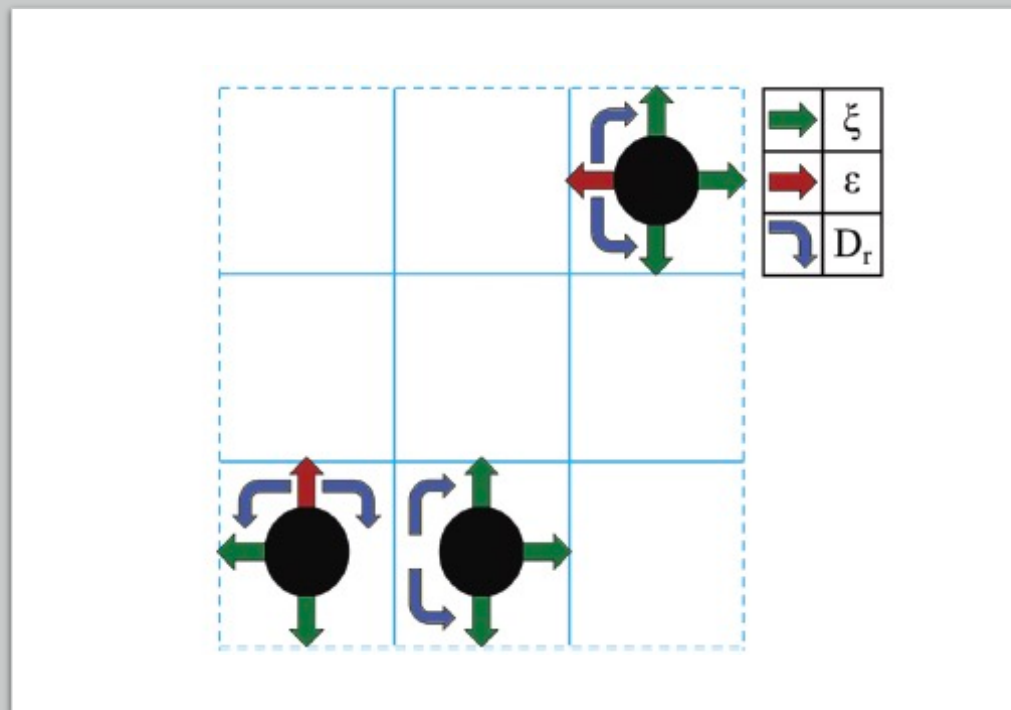


Ivo Buttinoni, Julian Bialké, Felix Kümmel, Hartmut Löwen, Clemens Bechinger, and Thomas Speck, Phys. Rev. Lett., 110, 238301 (2013)



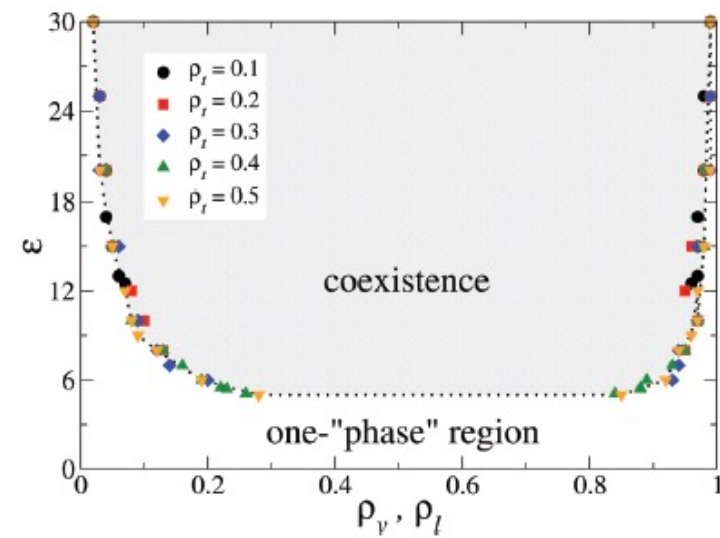
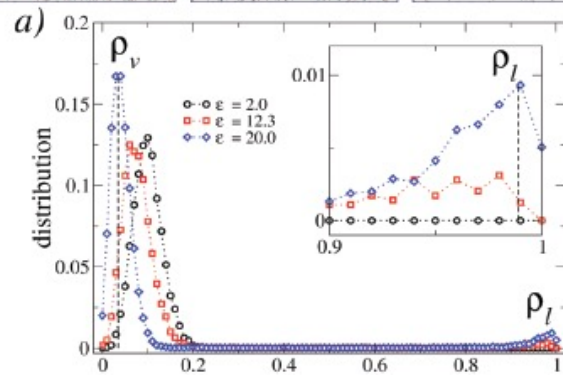
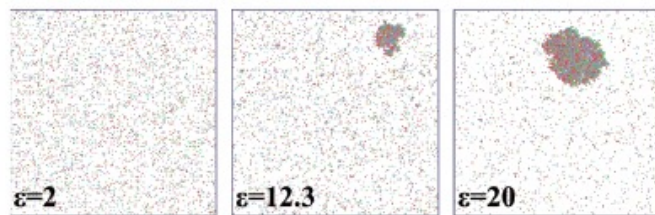
Wetting and interfacial roughness of Active Brownian Particles (ABP)

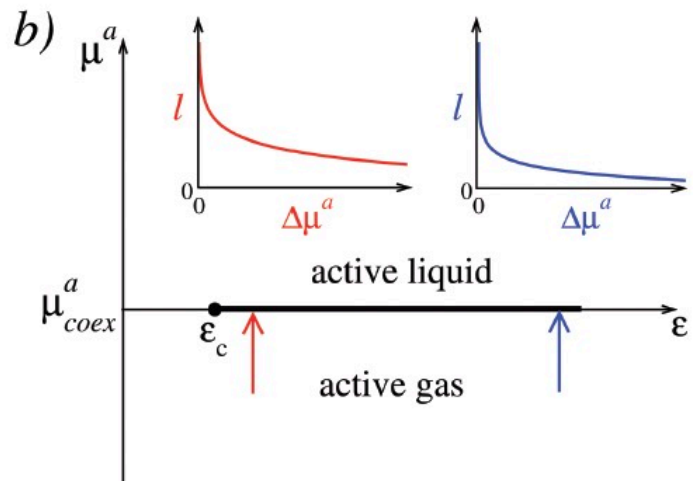
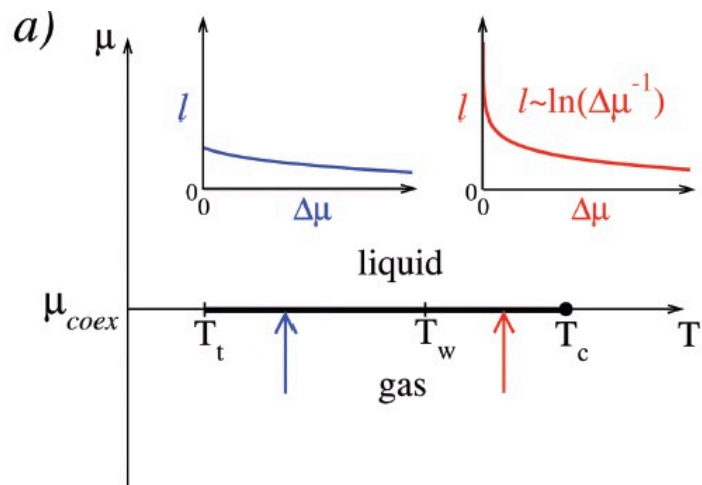
Lattice model of ABP: closed and open systems



Motility Induced Phase Separation (MIPS):

Activity vs Density Phase Diagram



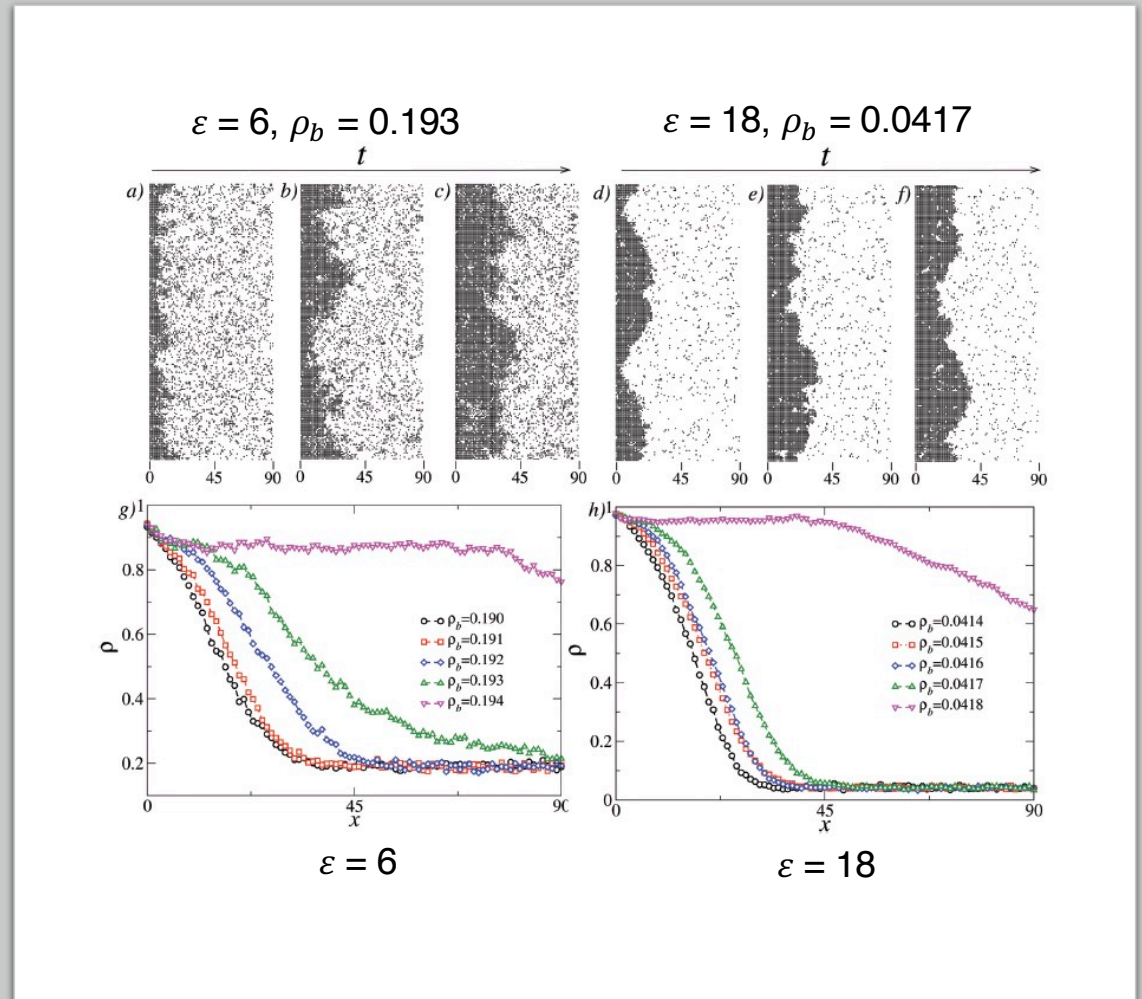


Recap

Complete wetting
phase diagram

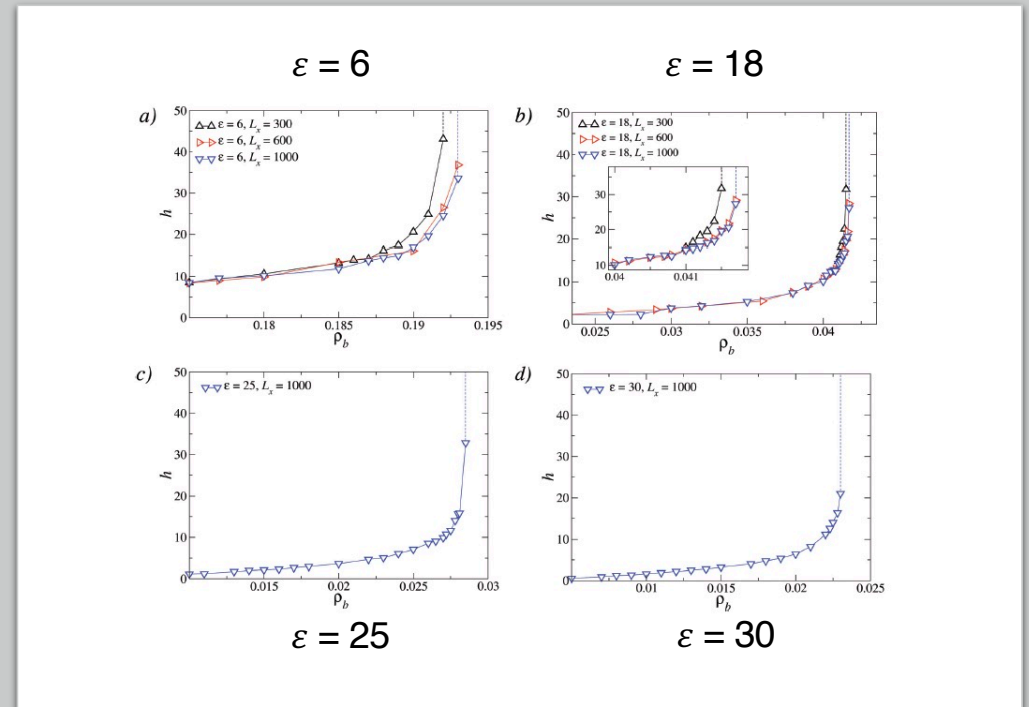
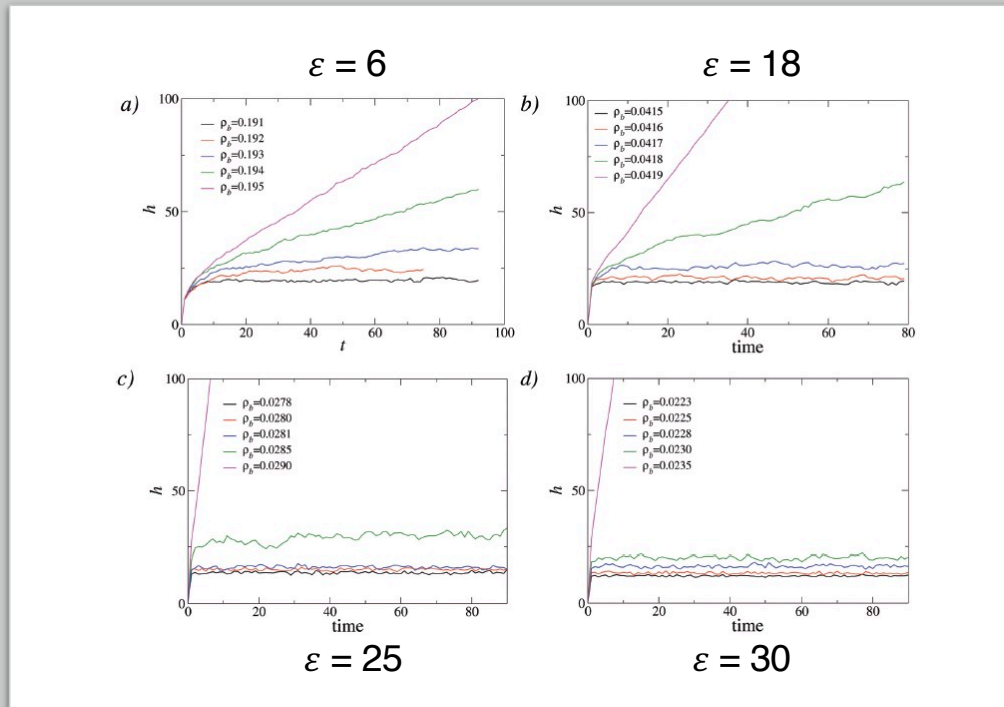
Density profiles:

different bulk densities at low and high activity

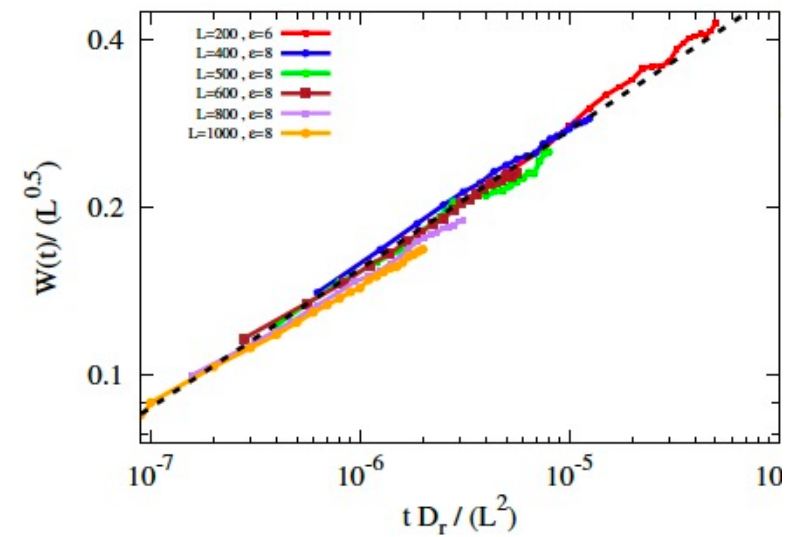
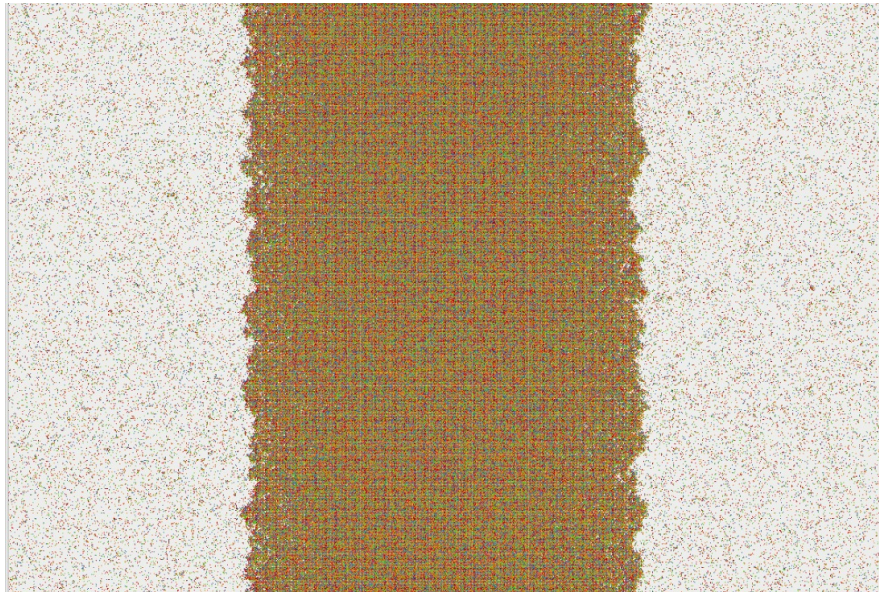


Film thickness:

different bulk densities at low and high activity

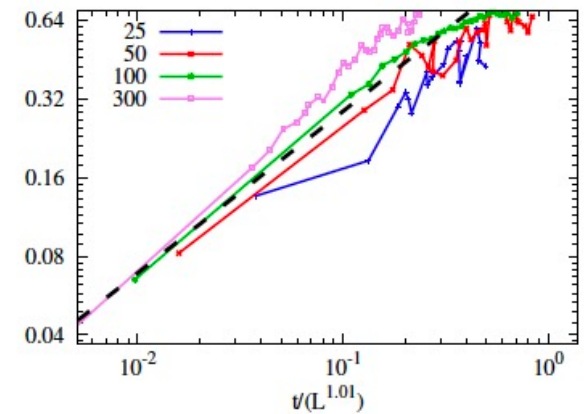
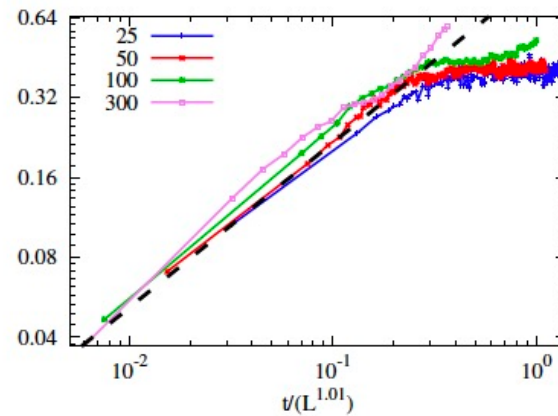
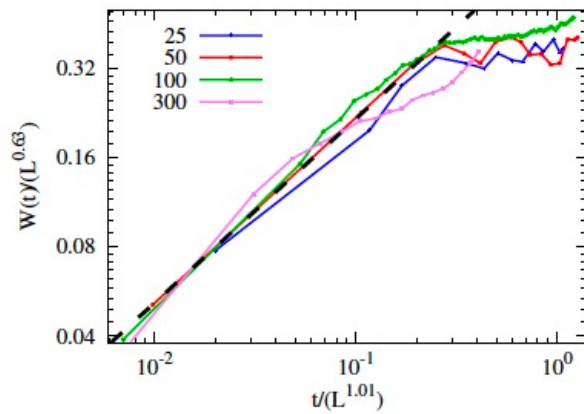


Roughness of the interface of a closed system: Edwards-Wilkinson



Roughness of the interface of an open system: KPZQ

Activity decreases and bulk density increases from left to right





Thanks

Pedro Neta, Cristóvão Dias, Mykola Tasinkevych and Nuno Araújo

